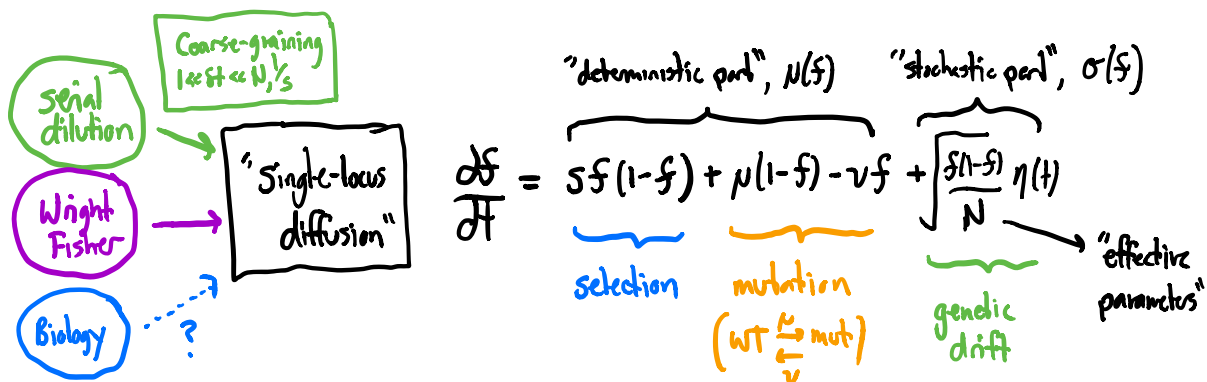


Announcements: Solutions for Problem Set 1 posted on Slack

Last time:



- Plan:
- ① can we understand this model mathematically? (-4 lectures)
 - ② Back to reality: DNA sequencing & genomics
 - ③ Multi-site genomes....

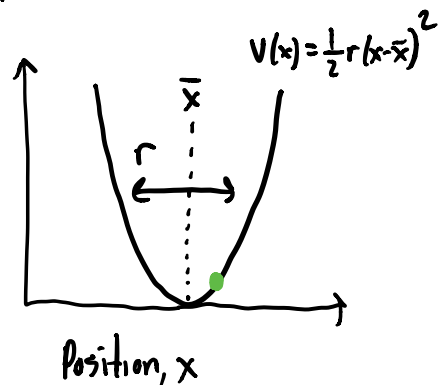
Last time: learn about SDEs w/ classic example:

"Brownian particle in a quadratic potential"

$$\frac{dx}{dt} = \underbrace{-r(x-\bar{x})}_{\text{restoring force}} + \underbrace{\sqrt{D}}_{\text{diffusion constant (}\propto kT \text{ in physics)}} \eta(t)$$

\bar{x}
equilibrium point

$-\frac{dV(x)}{dx}$



w/o noise ($D=0$), deterministic solution is $x_{det}(t) = x(0)e^{-rt} + \bar{x}(1 - e^{-rt})$
 (approaches \bar{x} @ rate r)

with noise?

\Rightarrow can focus on moments, e.g. mean $\langle x(t) \rangle \rightarrow N(0,1)$

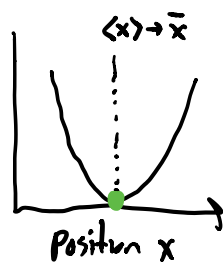
$$\begin{aligned} \langle x(t+\delta t) \rangle &= \langle x(t) - r[x(t) - \bar{x}]\delta t + \sqrt{D\delta t} Z_t \rangle \quad (\text{from "definition" of SDE}) \\ &= \langle x(t) \rangle - r[\langle x(t) \rangle - \bar{x}]\delta t + \sqrt{D\delta t} \langle Z_t \rangle \end{aligned}$$

When δt is small:

$$\Rightarrow \frac{\langle x(t+\delta t) - x(t) \rangle}{\delta t} \approx \frac{d\langle x(t) \rangle}{dt} = -r[\langle x(t) \rangle - \bar{x}] \quad (\text{ODE for } \langle x(t) \rangle)$$

\Rightarrow same as deterministic solution,

$$\langle x(t) \rangle = x(0)e^{-rt} + \bar{x}(1 - e^{-rt})$$



What about spread around this value?

\Rightarrow can look @ higher moments, e.g. if $\bar{x}=0$, want $\langle x(t)^2 \rangle$

can use same basic idea:

Step 1:

$$\langle x(t+\delta t)^2 \rangle = \langle [x(t) - r x(t) \delta t + \sqrt{D \delta t} z_t]^2 \rangle \quad \left(\begin{array}{l} \text{from "definition"} \\ \text{of } x(t+\delta t) \end{array} \right)$$

Step 2: expand to leading order in δt

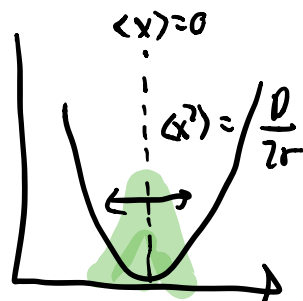
$$\begin{aligned} \langle x(t+\delta t)^2 \rangle &= \langle x(t)^2 - 2r x(t) \delta t \cdot x(t) + D \delta t z_t^2 + 2x \sqrt{D \delta t} z_t + \dots \rangle \\ &= \langle x(t)^2 \rangle - 2r \langle x(t)^2 \rangle \delta t + D \delta t \underbrace{\langle z_t^2 \rangle}_1 + 2 \langle x \rangle \sqrt{D \delta t} \underbrace{\langle z_t \rangle}_0 \end{aligned}$$

Step 3: take limit that δt small (reorganize)

$$\frac{\langle x(t+\delta t)^2 \rangle - \langle x(t)^2 \rangle}{\delta t} \approx \frac{\partial \langle x(t)^2 \rangle}{\partial t} = \underbrace{-2r \langle x^2 \rangle}_{\text{same as deterministic version}} + \underbrace{D}_{\substack{\uparrow \text{new part from} \\ \text{stochastic part.}}}$$

$$\Rightarrow \frac{\partial \langle x^2 \rangle}{\partial t} = 0 \Rightarrow \langle x^2 \rangle = \frac{D}{2r}$$

"balance between noise + deterministic restoring force.



can actually get full dist'n @ long times:

$$\frac{dx}{dt} = -\frac{dV(x)}{dx} + \sqrt{D} \eta(t) \Leftrightarrow \frac{dp(x,t)}{dt} = \frac{d}{dx} \left[-\frac{dV}{dx} p \right] + \frac{D}{2} \frac{d^2}{dx^2} [p]$$

("Fokker-Planck eq")

@ long times, $\frac{dp}{dt} = 0 \Rightarrow 0 = \frac{d}{dx} \left[-\frac{dV}{dx} p(x) \right] + \frac{D}{2} \frac{d^2 p}{dx^2}$

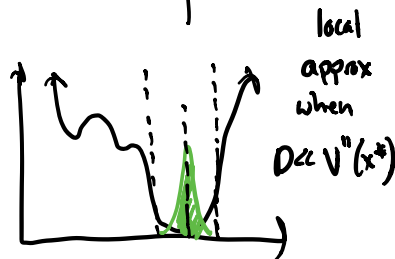
$$\Rightarrow -\frac{dV}{dx} p = \frac{D}{2} \frac{dp}{dx} \Rightarrow \frac{d \log p}{dx} = -\frac{2}{D} \frac{dV(x)}{dx}$$

$$\Rightarrow \log p(x) = -\frac{2V(x)}{D}$$

Gaussian dist'n
w/ mean \bar{x}
+
variance $\frac{D}{2r}$
↑

$$\Rightarrow p(x) \propto e^{-\frac{2V(x)}{D}} \propto e^{-\frac{1}{D} \frac{d^2 V(x^*)}{dx^2} (x-x^*)^2} \xrightarrow{\text{e.g.}} -\frac{c}{D} (x-\bar{x})^2$$

"Boltzmann distribution"



What about our evolutionary model?

e.g. $\frac{df}{dt} = sf(1-f) + \mu(1-f) - \nu f + \sqrt{\frac{f(1-f)}{N}} \eta(t)$ "Ito family stochastic process"

2 key differences: (1) $D_{\text{eff}} \approx \frac{f(1-f)}{N}$ (depends on mutation freq!)

(2) selection term ($sf(1-f)$) is non-linear

e.g. focus on $\langle f(t) \rangle$

Step 1:

$$\langle f(t + \delta t) \rangle = \left\langle f(t) + \delta t \left[sf(1-f) + \mu(1-f) - \nu f \right] + \sqrt{\frac{f(1-f)}{N}} \delta t Z_t \right\rangle$$

\uparrow
 $N(0,1)$

$$= \langle f(t) \rangle + \delta t \left[s \langle f \rangle - \langle f^2 \rangle + \mu(1 - \langle f \rangle) - \nu \langle f \rangle \right] + 0$$

$$\Rightarrow \frac{d\langle f(t) \rangle}{dt} = s \left[\langle f \rangle - \langle f^2 \rangle \right] + \mu(1 - \langle f \rangle) - \nu \langle f \rangle$$

\Downarrow
not $\langle f \rangle^2 \Rightarrow$ need $\langle f^2(t) \rangle$
to find $\langle f(t) \rangle$

do same thing for $\langle f(t+\delta t)^2 \rangle$

$$\dots \Rightarrow \frac{d\langle f^2 \rangle}{dt} = 2s \underbrace{\langle f \cdot f(1-f) \rangle}_{\text{from deterministic part}} + \underbrace{\frac{\langle f(1-f) \rangle}{N}}_{\text{from 2 stochastic terms, } \langle Z_i^2 \rangle = 1} + \text{mut'n.s.}$$

\Rightarrow depends on $\langle f^3 \rangle$ in addition to $\langle f \rangle$, $\langle f^2 \rangle$

\Rightarrow known as "moment hell" (general consequence of nonlinearity)

\Rightarrow one solution: focus on $s=0$ ("neutral theory")

What about the stationary distribution?

Fokker-Planck equation for evolutionary model:

$$\text{(remember: } \frac{dx}{dt} = \mu(x) + \sqrt{\sigma^2(x)} \eta(t) \Leftrightarrow \frac{dp}{dt} = -\frac{d}{dx}[\mu(x)p] + \frac{1}{2} \frac{d^2}{dx^2}[\sigma^2(x)p])$$

(1) here:

$$\frac{dp}{dt} = -\frac{d}{df} \left[[sf(1-f) + \mu(1-f) - \nu f] p \right] + \frac{1}{2} \frac{d^2}{df^2} \left(\frac{f(1-f)p}{N} \right)$$

\Rightarrow when $d_t p \approx 0$ (equilibrium)

$$p(f) \propto f^{-1} (1-f)^{-1} e^{-2N\Lambda(f)} \leftarrow$$

$$\text{where } \Lambda(f) = -[sf + \nu \log f + \nu \log(1-f)]$$

$$\Rightarrow \text{deterministically } \frac{d\Lambda}{df} = \frac{\partial \Lambda}{\partial f} \frac{df}{dt} = -\frac{1}{f(1-f)} \left(\frac{df}{dt} \right)^2 \leq 0$$

\Rightarrow dynamics try to "minimize" $\Lambda(f)$ [like "energy" $V(x)$]

$\Rightarrow N$ is analogy of " $\frac{1}{T_{\text{temp}}}$ " (or "noise")

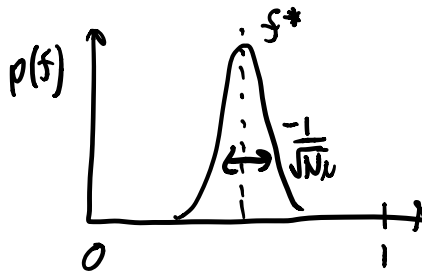
Plugging in for $\Lambda(f)$ in this example:

$$p(f) \propto f^{2N\mu-1} (1-f)^{2N\nu-1} e^{2Ns f} \quad \begin{array}{l} \text{"mutation-selection-drift} \\ \text{balance"} \\ \text{(Wright 1930s)} \end{array}$$

what does this look like?

Case 1: $N\mu, N\nu \gg 1$

peaked @ minimum of $\Lambda(f^*)$



$$\left. \frac{d\Lambda}{df} \right|_{f=f^*} = 0$$

"mutation-selection balance"

$$\Leftrightarrow s + \frac{\mu}{f^*} - \frac{\nu}{1-f^*} = 0$$

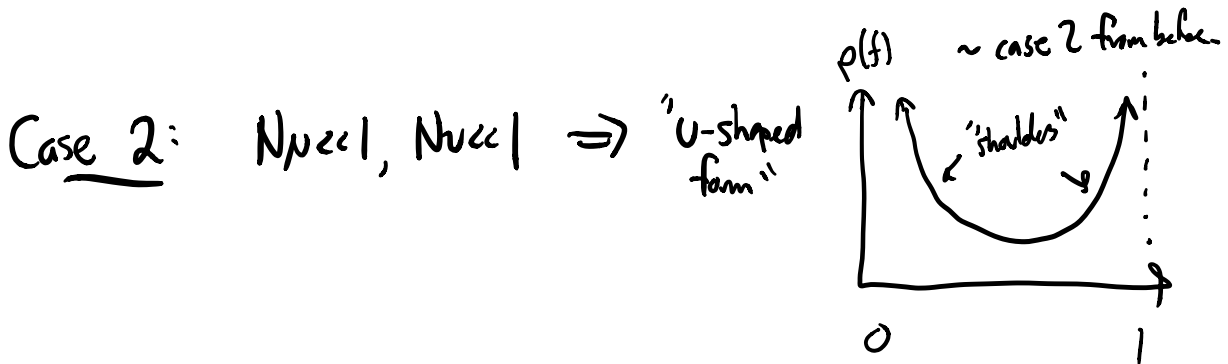
(same as deterministic sol'n,

$$\frac{df}{dt} = sf^*(1-f^*) + \mu(1-f^*) - \nu f^* = 0)$$

in simple limits:

$$f^* \approx \begin{cases} \frac{\mu}{|s|} & \text{if } s < 0; |s| \gg \mu, \nu \quad \begin{matrix} \downarrow s \\ \uparrow \mu \end{matrix} \\ \frac{\mu}{\mu + \nu} & \text{if } |s| \ll \mu, \nu \quad \begin{matrix} \downarrow \nu \\ \uparrow \mu \end{matrix} \\ 1 - \frac{\nu}{s} & \text{if } s \gg \mu, \nu \quad \begin{matrix} \downarrow \nu \\ \uparrow s \end{matrix} \end{cases}$$

full dist'n is $p(f) \propto (f^*)^{-1} (1-f^*)^{-1} e^{-N \frac{d^2\Lambda}{df^2}(f-f^*)^2}$



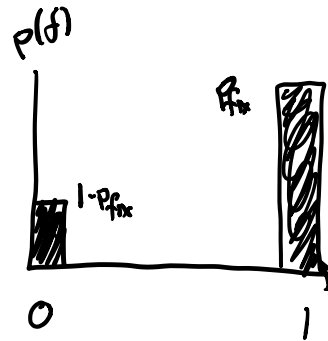
\Rightarrow when height of shoulders differs by $\sim e^{2Ns}$

\Rightarrow definitely not deterministic + a little noise.

$$\left(\frac{df}{dt} = sf(1-f) + \mu(1-f) - \nu f + \sqrt{\frac{f(1-f)}{N}} \eta(t) \right)$$

one mac stationary dist'n scenario:

no mutations: $\frac{df}{dt} = sf(1-f) + \sqrt{\frac{f(1-f)}{N}} \eta(t)$



$$p(f) = p_{\text{fix}} \delta(f-1) + (1-p_{\text{fix}}) \delta(f)$$

where p_{fix} depend on f_0

$$p_{\text{fix}}(f_0) \approx \begin{cases} 0 & \text{if } f_0=0 \\ ? & \\ 1 & \text{if } f_0=1 \end{cases}$$

Fokker-Planck Eq. not helpful, but generating function is!

$$H(z,t) \equiv \langle e^{-zf(t)} \rangle = \int e^{-zf} p(f,t) df$$

\Rightarrow use same approach...

$$\langle H(z,t+\delta t) \rangle = \langle e^{-z f(t+\delta t)} \rangle$$

$$= \left\langle e^{-z \left[f(t) + sf(1-f)\delta t + \sqrt{\frac{f(1-f)}{N}} \delta t z_t \right]} \right\rangle$$

= Taylor expand through $\mathcal{O}(\delta t)$ + avg over Z_t

$$= \underbrace{\langle e^{-z f(t)} \rangle}_{H(z,t)} + \left\langle e^{-z f} \left[\underbrace{-z s f(1-f)}_{\substack{\uparrow \\ \text{det} \\ \text{part}}} + \underbrace{\frac{z^2}{2N} f(1-f)}_{\substack{\uparrow \\ z \text{ stochastic terms.}}} \right] \delta t \right\rangle$$

\Rightarrow rearrange:

$$\frac{H(z,t+\delta t) - H(z,t)}{\delta t} = \frac{\partial H(z,t)}{\partial t} = \left\langle - \left[s z - \frac{z^2}{2N} \right] \underbrace{f(1-f)}_{\left[\left(-\frac{\partial}{\partial z} - \frac{\partial^2}{\partial z^2} \right) e^{-z f} \right]} e^{-z f} \right\rangle$$

$$\Rightarrow \frac{\partial H}{\partial t} = \left[s z - \frac{z^2}{2N} \right] \left[\frac{\partial H}{\partial z} + \frac{\partial^2 H}{\partial z^2} \right]$$

PDE for $H(z,t)$

hard in general, but for $z^* = 2Ns$

$$\frac{dH(z^*,t)}{dt} = 0 \Rightarrow H(z^*,t) = \text{const.} = e^{-z^* f_0}$$

$$H(z^*, \infty) = p_{\text{fix}} e^{-z^*} + (1-p_{\text{fix}}) e^{-0} = e^{-z^* f_0}$$

$$\Rightarrow P_{\text{fix}}(f_0) = \frac{1 - e^{-2Ns f_0}}{1 - e^{-2Ns}}$$

"fixation probability"
Kimura 1950's

$$\frac{df^2}{dt} = 2f \cdot \frac{df}{dt} = 2sf^2(1-f) + 2f \sqrt{\frac{f(1-f)}{N}} \eta(t)$$

$$\left\langle \frac{df^2}{dt} \right\rangle = \left\langle 2sf^2(1-f) \right\rangle + 0 \quad X$$

$$+ \left\langle \sqrt{\frac{f(1-f)}{N}} \cdot \sqrt{\frac{f(1-f)}{N}} \cdot \eta^2 \right\rangle \rightarrow 1$$

$$\frac{d\langle g(f) \rangle}{dt} = \left\langle \frac{dg}{df} \frac{df}{dt} \right\rangle + \left\langle \frac{1}{2} g''(f) \right\rangle \quad \text{"Ito formula"}$$